

Formation Sciences de la Matière Cours : Physique Master 1 (ENS)

Méthodes Mathématiques pour la Physique



TD 7

Thomas Buchert, CRAL, Observatoire de Lyon Rasha Boulos, ENS Clément Tauber, ENS thomas.buchert@ens-lyon.fr rasha.boulos@ens-lyon.fr clement.tauber@ens-lyon.fr

7.1. Geometric properties of the acceleration

Show that the acceleration is, in general, not orthogonal to the velocity. Show further that the acceleration always lies within the *osculating plane*, spanned by the unit tangent and unit normal vector fields of a differential curve.

7.2. Curvature of spacetime trajectories

A collection of masses subjected to a one-dimensional gravitational field strength trace curves in the two-dimensional spacetime. Employ the following equations and try to find their general solution giving a family of trajectories. Then, calculate their torsion and curvature.

The evolution of a (conserved) continuous density of masses ρ that move under a force field density, $\mathbf{f} = \rho \mathbf{g}$, is governed by Euler's equation, the continuity equation, and is subjected to a Newtonian gravitational field equation :

$$\frac{d}{dt}v_1 = g_1 \quad ; \tag{7.1}$$

$$\frac{d}{dt}\varrho = -\varrho\frac{\partial v_1}{\partial x_1} \quad ; \tag{7.2}$$

$$\frac{\partial g_1}{\partial x_1} = -4\pi\varrho \quad . \tag{7.3}$$

7.3. Averaged principal scalar invariants

Show that, by assuming existence of a velocity potential, $\mathbf{v} = \nabla S$, and by performing the spatial average over the principal scalar invariants of the velocity gradient field $(v_{i,j})$, we obtain :

$$\langle II \rangle_{\mathcal{D}} = \frac{1}{V_{\mathcal{D}}} \int_{\mathcal{D}} II(v_{i,j}) \ d^3x = \frac{1}{V_{\mathcal{D}}} \int_{\partial \mathcal{D}} H \ |\nabla S|^2 d\sigma \ ; \tag{7.4}$$

$$\langle III \rangle_{\mathcal{D}} = \frac{1}{V_{\mathcal{D}}} \int_{\mathcal{D}} III(v_{i,j}) \ d^3x = \pm \frac{1}{V_{\mathcal{D}}} \int_{\partial \mathcal{D}} G \left| \boldsymbol{\nabla} S \right|^3 d\sigma \quad , \tag{7.5}$$

where *H* is the local *mean curvature* and *G* the local *Gaussian curvature* at every point on the 2–surface bounding the domain. Recall that $|\nabla S| = \frac{ds}{dt}$ with the instrinsic arc length *s* of the trajectories of fluid elements, and the extrinsic Newtonian time *t*.

Hint : Make explicit use of the properties of gradient fields and geometrical properties of surfaces appended below.

$$I(v_{i,j}) = \Theta = \nabla \cdot \mathbf{v} ; \quad 2II(v_{i,j}) = \nabla \cdot \left(\mathbf{v} (\nabla \cdot \mathbf{v}) - (\mathbf{v} \cdot \nabla) \mathbf{v} \right) ; \quad 3III(v_{i,j}) = \nabla \cdot \left(\frac{1}{2} \nabla \cdot \left(\mathbf{v} (\nabla \cdot \mathbf{v}) - (\mathbf{v} \cdot \nabla) \mathbf{v} \right) \mathbf{v} - \left(\mathbf{v} (\nabla \cdot \mathbf{v}) - (\mathbf{v} \cdot \nabla) \mathbf{v} \right) \cdot \nabla \mathbf{v} \right).$$

$$(7.6)$$

At points $P = (x_0, y_0, z_0)$, where the representation of the velocity front in terms of surfaces in the form $z = \chi(x, y)$ is nonsingular, $\nabla S \neq 0$, we have for the mean curvature H and the Gaussian curvature G (Recall that indexed letters denote partial derivatives with respect to the coordinates) :

$$2H := \frac{(1+\chi_x^2)\chi_{yy} - 2\chi_x\chi_y\chi_{xy} + (1+\chi_y^2)\chi_{xx}}{(1+\chi_x^2 + \chi_y^2)^{3/2}} ;$$
(7.7)

$$G := \frac{\chi_{xx}\chi_{yy} - \chi_{xy}^2}{(1 + \chi_x^2 + \chi_y^2)^2}$$
(7.8)

Using the implicit definition $S(x, y, \chi(x, y)) = s$ of the velocity front, calculate the derivatives of χ to obtain for the curvature invariants of the front :

$$2H = \frac{1}{|\nabla S|^3} \left[2S_x S_y S_{xy} + 2S_x S_z S_{xz} + 2S_y S_z S_{yz} - S_{xx} (S_y^2 + S_z^2) - S_{yy} (S_x^2 + S_z^2) - S_{zz} (S_x^2 + S_y^2) \right] ;$$

$$G = \frac{1}{|\nabla S|^4} \left[S_x^2 (S_{yy} S_{zz} - S_{yz}^2) + S_y^2 (S_{xx} S_{zz} - S_{xz}^2) + S_z^2 (S_{xx} S_{yy} - S_{xy}^2) - 2S_x S_y (S_{xy} S_{zz} - S_{xz} S_{yz}) - 2S_x S_z (S_{xz} S_{yy} - S_{xy} S_{yz}) - 2S_y S_z (S_{yz} S_{xx} - S_{xy} S_{xz}) \right] .$$

$$(7.9)$$